0.1 The Eigenvalue Problem

$$\frac{d^2\Psi(x^*)}{dx^{*2}} - l^2\Psi(x^*) = 0, \quad \Psi(0) = 0, \quad \Sigma(1) = 0,
\frac{d^2\Sigma(x^*)}{dx^{*2}} - l^2\Sigma(x^*) = 0, \quad \Psi(S_0^*) = \Sigma(S_0^*) + \Phi N_0, ,$$
(1)

where

$$N_0 = F_1(S_0^*) \frac{C}{k} \frac{Pec_i}{R} \left\{ (R-1) + \frac{k R R_3 (1-R_1)}{C Pec_i} \right\}.$$

$$F(S_0^*) = \frac{\dot{m}_{liq,vap}}{\dot{m}_i} = F_1(S_0^*) \left[1 - R G_1 S_0^* \right], \quad F_1(S_0^*) = \frac{1}{R S_0^* + 1 - S_0^*},$$

$$G_1 = \frac{R_3 k (1 - R_1)}{C Pec_i}.$$
(2)

The eigenvalue problem for the temperature distributions in the both phases,

$$\left(\frac{d^{2}}{dx^{*2}} - \omega_{0} \frac{d}{dx^{*}} - E_{1} \sigma - l^{2}\right) \phi_{liq} + \left. \frac{d\Theta_{liq}^{0}}{dx^{*}} \frac{d\Psi}{dx^{*}} = 0,$$

$$\phi_{liq}(0) = 0, \quad \phi_{liq}(S_{0}^{*}) = -\Phi \left. \frac{d\Theta_{liq}^{0}}{dx^{*}} \right|_{x^{*} = S_{0}^{*}},$$

$$\left(\frac{d^{2}}{dx^{*2}} - \Omega_{0} \frac{d}{dx^{*}} - \frac{E_{2} k R_{1}}{C} \sigma - l^{2}\right) \phi_{vap} + \frac{R k}{C} \frac{d\Theta_{vap}^{0}}{dx^{*}} \frac{d\Sigma}{dx^{*}} = 0,$$

$$\phi_{vap}(1) = 0, \quad \phi_{vap}(S_{0}^{*}) = -\Phi \left. \frac{d\Theta_{vap}^{0}}{dx^{*}} \right|_{x^{*} = S^{*}},$$
(3)

where

$$\frac{d\Theta_{liq}^{0}}{dx^{*}} = N_{1} \exp(\omega_{0} x^{*}),$$

$$\frac{d\Theta_{vap}^{0}}{dx^{*}} = N_{2} \exp(\Omega_{0} (x^{*} - 1)),$$
(4)

with leads to

$$\frac{d\Theta_{liq}^{0}}{dx^{*}}\Big|_{x^{*}=S_{0}^{*}} = N_{3}, \text{ and } \frac{d\Theta_{vap}^{0}}{dx^{*}}\Big|_{x^{*}=S_{0}^{*}} = N_{4},$$
 (5)

where

$$N_{1} = \frac{\omega_{0}}{\exp(\omega_{0} S_{0}^{*}) - 1}, \quad N_{2} = \frac{\Omega_{0}}{\exp(\Omega_{0} (S_{0}^{*} - 1)) - 1},$$

$$N_{3} = N_{1} \exp(\omega_{0} S_{0}^{*}), \quad N_{4} = N_{2} \exp(\Omega_{0} (S_{0}^{*} - 1)).$$

$$(6)$$

and

$$\omega_0 = \frac{1}{k} \operatorname{Pec}_i F(S_0^*) C,$$

$$\Omega_0 = \operatorname{Pec}_i F(S_0^*),$$
(7)

0.2 Dispersion Analysis

An equation which represent the relationship between the decay rate σ and the wave number l is known as dispersion equation and can be obtained by using the energy jump across the liquid-vapour interface.

$$\varphi \, \sigma \, H_{liq} \, \Phi = \left\{ \Phi \, \frac{d^2 \Theta_{liq}^0}{dx^{*\,2}} + \frac{d\phi_{liq}}{dx^*} \right\} + \frac{\Theta_0}{k} \left\{ \Phi \, \frac{d^2 \Theta_{vap}^0}{dx^{*\,2}} + \frac{d\phi_{vap}}{dx^*} \right\} - H_{liq} \, \frac{d\Psi_1}{dx^*}. \tag{8}$$

From the basic state we know that

$$\frac{d^{2}\Theta_{liq}^{0}}{dx^{*2}}\Big|_{x^{*}=S_{0}^{*}} = N_{3}\omega_{0},$$

$$\frac{d^{2}\Theta_{vap}^{0}}{dx^{*2}}\Big|_{x^{*}=S_{0}^{*}} = N_{4}\Omega_{0}.$$
(9)

Substituting (9) into (8) yields

$$\varphi \,\sigma \,H_{liq} \,\Phi = \left\{ \frac{d\phi_{liq}}{dx^*} + \Phi \,N_3 \,\omega_0 \right\} + \frac{\Theta_0}{k} \left\{ \frac{d\phi_{vap}}{dx^*} + \Phi \,N_4 \,\Omega_0 \right\} - H_{liq} \,\frac{d\Psi_1}{dx^*}. \tag{10}$$

The complete solution for $\sigma(l)$ can be obtained by substituting ϕ_{liq} , ϕ_{vap} and Ψ_1 into (10) which is solved numerically using Maple.