

My Question : please variation this action with respect to tensor metric $g_{\mu\nu}$.this called the Einstein equation. To obtain the Einstein equation, we vary the action with respect to the metric tensor

$$s = \int d^4x \sqrt{-g} [R + \alpha(\varphi) R_{GB}]$$

$$R_{GB} = R^2 - 4R_{\mu\nu}R^{\mu\nu} + R_{\mu\nu\rho\sigma}R^{\mu\nu\rho\sigma}$$

where R is the Ricci scalar, $R_{\{\mu\nu\}}$ is a Ricci tensor, $R_{\{\mu\nu\rho\sigma\}}$ is a Riemann tensor, g is the determinant of the metric tensor, $\alpha(\varphi)$ is a function of $\varphi(r)$, R_{GB} is a quadratic Gauss-bonnet (GB) term, G is the gravitational constant, and (d^4x) represents the volume element in four-dimensional spacetime.

$$\alpha(\varphi) = f(\varphi)$$

$$\varphi = \varphi(r)$$

Let us now turn our attention to the third term δS_3 in (19.34), in which we must express $\delta\sqrt{-g}$ in terms of the variation $\delta g^{\mu\nu}$. Recalling that $g = \det[g_{\mu\nu}]$, we note that the cofactor of the element $g_{\mu\nu}$ in this **determinant** is $gg_{\mu\nu}$. It follows that

$$\delta g = gg^{\mu\nu} \delta g_{\mu\nu} = -gg_{\mu\nu} \delta g^{\mu\nu},$$

where in the second equality we have used the result (19.33). Thus, we have

$$\delta\sqrt{-g} = -\frac{1}{2}(-g)^{-1/2} \delta g = -\frac{1}{2}\sqrt{-g} g_{\mu\nu} \delta g^{\mu\nu}.$$

I see different results in papers, but I don't have any code in Maple for them. I need the Maple code for this equation to continue deriving my equation.

so i give you some example of this result in next page :

Example 1: (arXiv:1812.06941v1 [hep-th] 17 Dec 2018)

$$S = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left[R - \frac{1}{2} \partial_\mu \phi \partial^\mu \phi + f(\phi) R_{GB}^2 - 2\Lambda \right].$$

$$R_{GB}^2 = R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma} - 4R_{\mu\nu} R^{\mu\nu} + R^2,$$

by varying the action (S) with respect to the metric tensor $g_{\mu\nu}$, we derive the gravitational field equations:

$$G_{\mu\nu} = T_{\mu\nu}$$

where $G_{\mu\nu}$ is the Einstein tensor and $T_{\mu\nu}$ is the energy-momentum tensor, with the latter having the form

$$T_{\mu\nu} = -\frac{1}{4} g_{\mu\nu} \partial_\rho \phi \partial^\rho \phi + \frac{1}{2} \partial_\mu \phi \partial_\nu \phi - \frac{1}{2} (g_{\rho\mu} g_{\lambda\nu} + g_{\lambda\mu} g_{\rho\nu}) \eta^{\kappa\lambda\alpha\beta} \tilde{R}^{\rho\gamma}{}_{\alpha\beta} \nabla_\gamma \partial_\kappa f(\phi) - \Lambda g_{\mu\nu}.$$

example2 : (arXiv:hep-th/0504052v2 25 May 2005)

The starting action is given by

$$S = \int d^4x \sqrt{-g} \left\{ \frac{1}{2\kappa^2} R - \frac{\gamma}{2} \partial_\mu \phi \partial^\mu \phi - V(\phi) + f(\phi) G \right\}.$$

On the other hand, the variation over the metric $g_{\mu\nu}$ gives

$$\begin{aligned} 0 = & \frac{1}{\kappa^2} \left(-R^{\mu\nu} + \frac{1}{2} g^{\mu\nu} R \right) + \gamma \left(\frac{1}{2} \partial^\mu \phi \partial^\nu \phi - \frac{1}{4} g^{\mu\nu} \partial_\rho \phi \partial^\rho \phi \right) + \frac{1}{2} g^{\mu\nu} (-V(\phi) + f(\phi) G) \\ & - 2f(\phi) R R^{\mu\nu} + 2\nabla^\mu \nabla^\nu (f(\phi) R) - 2g^{\mu\nu} \nabla^2 (f(\phi) R) \\ & + 8f(\phi) R^\mu{}_\rho R^{\nu\rho} - 4\nabla_\rho \nabla^\mu (f(\phi) R^{\nu\rho}) - 4\nabla_\rho \nabla^\nu (f(\phi) R^{\mu\rho}) \\ & + 4\nabla^2 (f(\phi) R^{\mu\nu}) + 4g^{\mu\nu} \nabla_\rho \nabla_\sigma (f(\phi) R^{\rho\sigma}) - 2f(\phi) R^{\mu\rho\sigma\tau} R^\nu{}_{\rho\sigma\tau} + 4\nabla_\rho \nabla_\sigma (f(\phi) R^{\mu\rho\sigma\nu}). \end{aligned}$$